

# **Sgoldstino interpretation of HyperCP events**

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# Outline

- $\Sigma^+$  decay in HyperCP experiment
- Sgoldstino, what is it?
- Decays of  $D^-$  and  $B^-$ -mesons
- Decays of  $\rho^-$ ,  $\omega^-$  and  $\phi^-$ -mesons
- Conclusions

# HyperCP events

- Three events of  $\Sigma^+ \rightarrow p\mu^+\mu^-$  (HyperCP, Fermilab, 2005)
- $\text{Br}(\Sigma^+ \rightarrow p\mu^+\mu^-) = [3.1_{-1.9}^{+2.4}(\text{stat.}) \pm 1.5(\text{sys.})] \times 10^{-8}$
- $\Sigma \rightarrow pX$ , where  $X \rightarrow \mu^+\mu^-$  with  $m_X = 214.3 \pm 0.5$  MeV
- $X$  can be pseudoscalar or pseudovector, but not scalar or vector
- $X$  — pseudoscalar sgoldstino

# Sgoldstino — what is it?

- SUSY theories
- Superfields:  $\hat{S} = \phi + \sqrt{2}\theta\psi + \theta^2\tilde{F}$
- $\langle\tilde{F}\rangle = F$  — SUSY broken spontaneously
  - $\sqrt{F}$  — the scale of SUSY breaking
- Supersymmetric analog of Nambu-Goldstone theorem:  
 $\psi$  is massless fermion — goldstino
- Sgoldstino:  $\phi = \frac{1}{\sqrt{2}}(S + iP)$ :  
 $S$  — scalar and  $P$  — pseudoscalar

# Interactions of sgoldstino

- Couplings to the Standard model fields are related with soft SUSY breaking terms

An example:

SUSY breaking lagrangian

$$\mathcal{L}_{soft} = -m_{H_u}^2 |H_u|^2 - m_{H_d}^2 |H_d|^2 - B\mu H_u H_d$$

can be obtained from

$$\begin{aligned} \mathcal{L}_{SUSY} = \int d^2\theta d^2\bar{\theta} & \left( -\frac{m_{H_u}^2}{F} \hat{S} \hat{H}_u^\dagger \hat{H}_u - \frac{m_{H_d}^2}{F} \hat{S} \hat{H}_d^\dagger \hat{H}_d \right) \\ & + \int d^2\theta \left( -\frac{B\mu}{F} \hat{S} \hat{H}_u \hat{H}_d \right) \end{aligned}$$

- Masses  $m_S$  and  $m_P$  are not related with the soft terms
- Sgoldstinos  $S$  and  $P$  can be light

# Sgoldstino in $\Sigma^+$ decay

- Lagrangian for  $s - d$  transition:

$$\mathcal{L}_{Pds} = -P \cdot (h_{12}^{(D)} \cdot \bar{d} i\gamma^5 s + \text{h.c.}) ,$$

- From measured branching one obtains:

$$|h_{12}^{(D)}| \cdot \text{Br}^{1/2}(P \rightarrow \mu^+ \mu^-) = 3.8 \cdot 10^{-10}$$

- Constraints on sgoldstino lifetime:

$$1.7 \cdot 10^{-15} \text{ s} \lesssim \tau_P \lesssim 2.5 \cdot 10^{-11} \text{ s}$$

- The energy scale of SUSY breaking is low :

$$2 \text{ TeV} \lesssim \sqrt{F} \lesssim 65 \text{ TeV}$$

- $m_P = 214.3 \text{ MeV}$

- Possible decays:

$$P \rightarrow \mu^+ \mu^-, P \rightarrow e^+ e^- \text{ and } P \rightarrow \gamma\gamma$$

$$\frac{\Gamma(P \rightarrow e^+ e^-)}{\Gamma(P \rightarrow \mu^+ \mu^-)} \sim \frac{m_e^2}{m_\mu^2}, \quad \frac{\Gamma(P \rightarrow \gamma\gamma)}{\Gamma(P \rightarrow \mu^+ \mu^-)} \sim 1 \div 10^4$$

- Predicted branchings of kaon decays with  $s - d$  transition:

$$\text{Br}(K^+ \rightarrow \pi^+ \pi^0 P(P \rightarrow \mu^+ \mu^-)) \sim 10^{-12}$$

$$\text{Br}(K_L \rightarrow \pi^0 \pi^0 P(P \rightarrow \mu^+ \mu^-)) \sim 10^{-8}$$

$$\text{Br}(K_S \rightarrow \pi^0 \pi^0 P(P \rightarrow \mu^+ \mu^-)) \sim 10^{-11}$$

# Sgoldstino interactions with quarks

- General interaction lagrangian:

$$\mathcal{L}_P = -P \cdot \left( h_{jl}^{(U)} \cdot \bar{u}_j i\gamma^5 u_l + h_{jl}^{(D)} \cdot \bar{d}_j i\gamma^5 d_l + \text{h.c.} \right),$$

- Constants  $h_{jl} = \frac{m_{jl}^{(LR)2}}{\sqrt{2}F}$  depend on model

- Three models:

- I:  $h_{jl} \sim h_{12}^D$ .
- II:  $h_{jl} \sim \frac{A}{F} \max(m_j, m_l)$
- III: particular left-right SUSY model

# Decays of $D$ - and $B$ -mesons

- Interactions with quarks induce flavor changing neutral transitions
- Possible signature of sgoldstino:

$$P_{B,D} \rightarrow \mathcal{V} P ,$$

- For decay with  $u_l - u_j$  or  $d_l - d_j$  transition:

$$\mathcal{M} (P_{B,D} \rightarrow \mathcal{V} P) = h_{jl} \cdot P(q) \cdot \mathcal{O}_{P_{B,D},\mathcal{V}} (q^2) ,$$

where  $q$  — sgoldstino momentum and

$$\mathcal{O}_{P_{B,D},\mathcal{V}} (q^2) \equiv \langle \mathcal{V} (p_{\mathcal{V}}, \varepsilon_{\mathcal{V}}) | \bar{q}_j \gamma^5 q_l | P_{B,D}(p) \rangle =$$
$$A_0^{(P_{B,D},\mathcal{V})} (q^2) \cdot \frac{-2im_{\mathcal{V}}}{m_j + m_l}$$

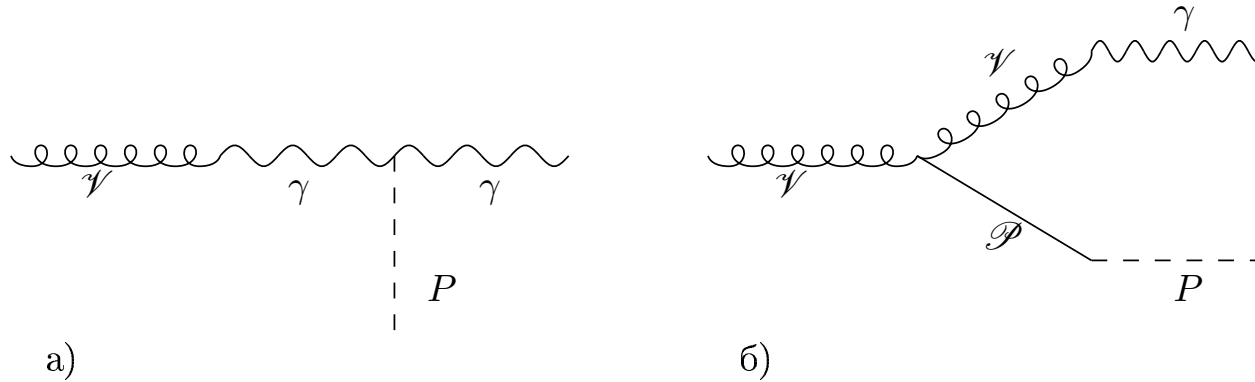
# Numerical results

Decay	$h_{jl}$	$\text{Br}_{(I)}$	$\text{Br}_{(II)}$	$\text{Br}_{(III)}$
$B_s \rightarrow \phi P(P \rightarrow \mu^+ \mu^-)$	$h_{23}^{(D)}$	$6.5 \cdot 10^{-9}$	$8.8 \cdot 10^{-6}$	$8.7 \cdot 10^{-6}$
$B_s \rightarrow K^{*0} P(P \rightarrow \mu^+ \mu^-)$	$h_{13}^{(D)}$	$5.3 \cdot 10^{-9}$	$7.2 \cdot 10^{-6}$	$2.3 \cdot 10^{-7}$
$B_c^+ \rightarrow D^{*+} P(P \rightarrow \mu^+ \mu^-)$	$h_{13}^{(D)}$	$3.2 \cdot 10^{-10}$	$4.4 \cdot 10^{-7}$	$1.4 \cdot 10^{-8}$
$B_c^+ \rightarrow D_s^{*+} P(P \rightarrow \mu^+ \mu^-)$	$h_{23}^{(D)}$	$3.0 \cdot 10^{-10}$	$4.0 \cdot 10^{-7}$	$4.0 \cdot 10^{-7}$
$B_c^+ \rightarrow B^{*+} P(P \rightarrow \mu^+ \mu^-)$	$h_{12}^{(U)}$	$4.1 \cdot 10^{-10}$	$4.4 \cdot 10^{-8}$	$8.2 \cdot 10^{-7}$
$B^+ \rightarrow K^{*+} P(P \rightarrow \mu^+ \mu^-)$	$h_{23}^{(D)}$	$3.8 \cdot 10^{-9}$	$5.2 \cdot 10^{-6}$	$5.1 \cdot 10^{-6}$
$B^0 \rightarrow K^{*0} P(P \rightarrow \mu^+ \mu^-)$		$3.5 \cdot 10^{-9}$	$4.8 \cdot 10^{-6}$	$4.7 \cdot 10^{-6}$
$B^0 \rightarrow \rho P(P \rightarrow \mu^+ \mu^-)$	$h_{13}^{(D)}$	$3.1 \cdot 10^{-9}$	$4.2 \cdot 10^{-6}$	$1.4 \cdot 10^{-7}$
$B^+ \rightarrow \rho^+ P(P \rightarrow \mu^+ \mu^-)$		$3.3 \cdot 10^{-9}$	$4.6 \cdot 10^{-6}$	$1.3 \cdot 10^{-7}$
$D^0 \rightarrow \rho P(P \rightarrow \mu^+ \mu^-)$	$h_{12}^{(U)}$	$1.4 \cdot 10^{-9}$	$1.5 \cdot 10^{-7}$	$2.8 \cdot 10^{-6}$
$D^+ \rightarrow \rho^+ P(P \rightarrow \mu^+ \mu^-)$		$3.5 \cdot 10^{-9}$	$3.7 \cdot 10^{-7}$	$7.0 \cdot 10^{-6}$

$$\text{Br}(B^+ \rightarrow K^{*+} \mu^+ \mu^-) < 2.2 \times 10^{-6},$$

# Decays of $\rho$ -, $\omega$ - and $\phi$ -mesons

- Another possible signature of sgoldstino:  $\psi \rightarrow P\gamma$ .
- Two contributions to matrix element



$$\mathcal{L}_1 = \frac{1}{4\sqrt{2}} \varepsilon^{\mu\nu\lambda\rho} \left( \frac{M_{\gamma\gamma}}{F} P F_{\mu\nu} F_{\lambda\rho} + \frac{M_3}{F} P G_{\mu\nu}^a G_{\lambda\rho}^a \right),$$

where  $M_{\gamma\gamma} = M_1 \cos^2 \theta_w + M_2 \sin^2 \theta_w$

$$\mathcal{L}_{VMD} = \frac{e}{g} A_\mu \left( \sqrt{2} m_\rho^2 \rho^\mu + \frac{\sqrt{2}}{3} m_\omega^2 \omega^\mu - \frac{2}{3} m_\phi \phi^\mu \right), \quad g \simeq 8.6$$

# Decays of $\rho$ -, $\omega$ - and $\phi$ -mesons

- Mixing of sgoldstino with  $\pi^0$ ,  $\eta$  and  $\eta'$

$$\mathcal{L}_{mix} = C_{P\pi^0} P \pi^0 + C_{P\eta} P \eta + C_{P\eta'} P \eta'$$

- Interactions of vector and pseudoscalar mesons with photon

$$\mathcal{L}_{\mathcal{V}\mathcal{P}\gamma} = e g_{\mathcal{V}\mathcal{P}\gamma} \epsilon_{\mu\nu\lambda\rho} \partial_\mu A_\nu \partial_\lambda \mathcal{V}_\rho \mathcal{P}.$$

where  $g_{\mathcal{V}\mathcal{P}\gamma}$  are tuned to saturate the decays  $\mathcal{P} \rightarrow \mathcal{V} \gamma$  or  $\mathcal{V} \rightarrow \mathcal{P} \gamma$ .

# Numerical results

- The width of the decay:

$$\Gamma(\psi \rightarrow P\gamma) = \frac{C_\psi^2}{96\pi} \frac{(m_\psi^2 - m_P^2)^3}{m_\psi^3},$$

where

$$C_\psi = -\frac{\sqrt{2}a_\psi eM_{\psi\gamma}}{gF} + \sum_{\mathcal{P}=\pi^0,\eta,\eta'} \frac{eg_\psi \mathcal{P}_\gamma C_{P\mathcal{P}}}{m_P^2 - m_{\mathcal{P}}^2}.$$

$$1.8 \times 10^{-13} < Br(\phi \rightarrow \gamma P(P \rightarrow \mu^+ \mu^-)) < 1.6 \times 10^{-7},$$

$$9.1 \times 10^{-15} < Br(\rho \rightarrow \gamma P(P \rightarrow \mu^+ \mu^-)) < 3.3 \times 10^{-7},$$

$$1.8 \times 10^{-14} < Br(\omega \rightarrow \gamma P(P \rightarrow \mu^+ \mu^-)) < 3.4 \times 10^{-7}.$$

# Conclusions

- We considered the possibility to check the sgoldstino interpretation of the HyperCP events in the decays of  $D^-$ ,  $B^-$ ,  $\rho^-$ ,  $\omega^-$  and  $\phi^-$ -mesons
- Estimated branchings of the decays  $P_{B,D} \rightarrow VP (P \rightarrow \mu^+ \mu^-)$  (and  $m_P = 214$  MeV) can be at the level  $10^{-8} \div 10^{-6}$
- The branchings of the decays  $\rho^-$ ,  $\omega^-$  and  $\phi^-$ -mesons  $\psi \rightarrow P\gamma$  are in the range  $10^{-13} \div 10^{-7}$