



# Слияние нейтронных звёзд и гамма-всплески: модель обдирания

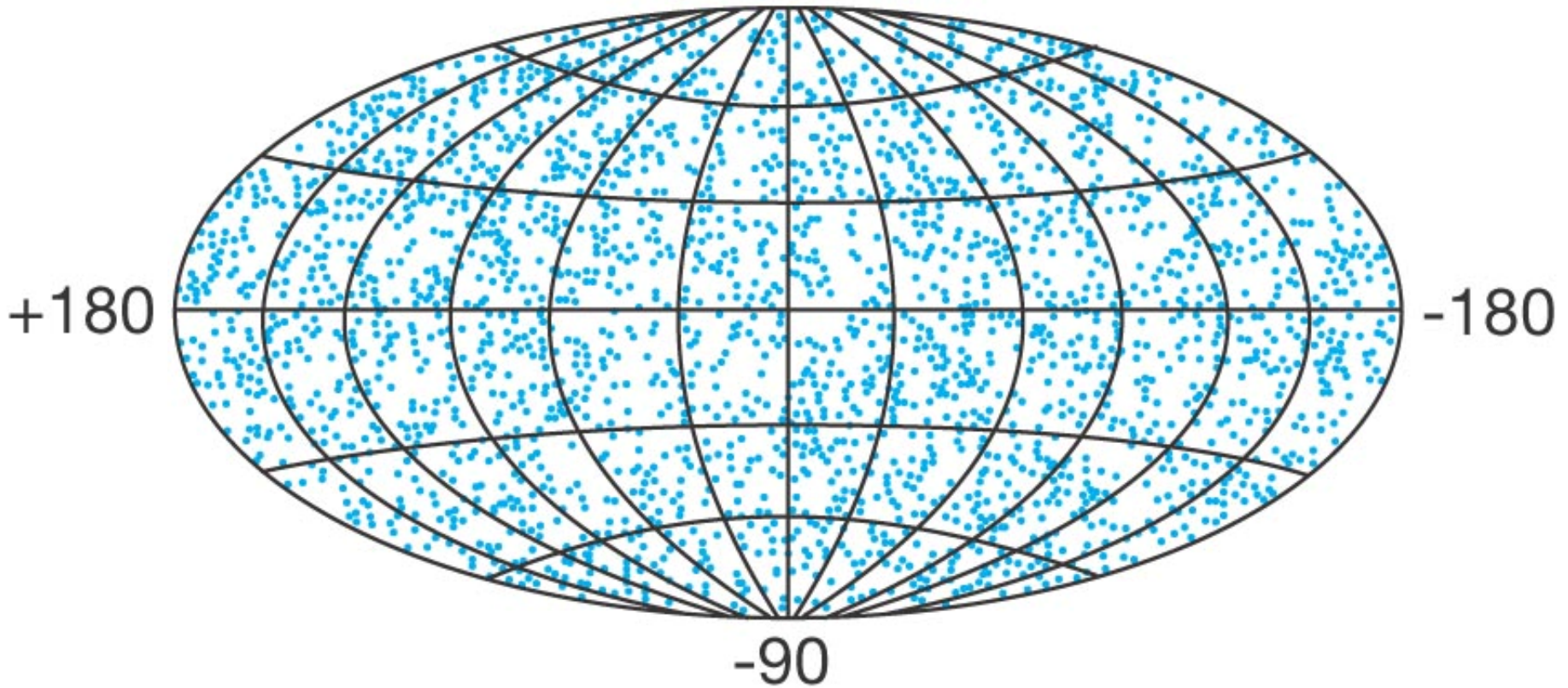
**С.И. Блинников, Д.К. Надёжин и А.В. Юдин**

Десятые Зацепинские чтения, 7 июня 2019 г.

# Gamma-Ray Bursts

2704 gamma-ray bursts

+90



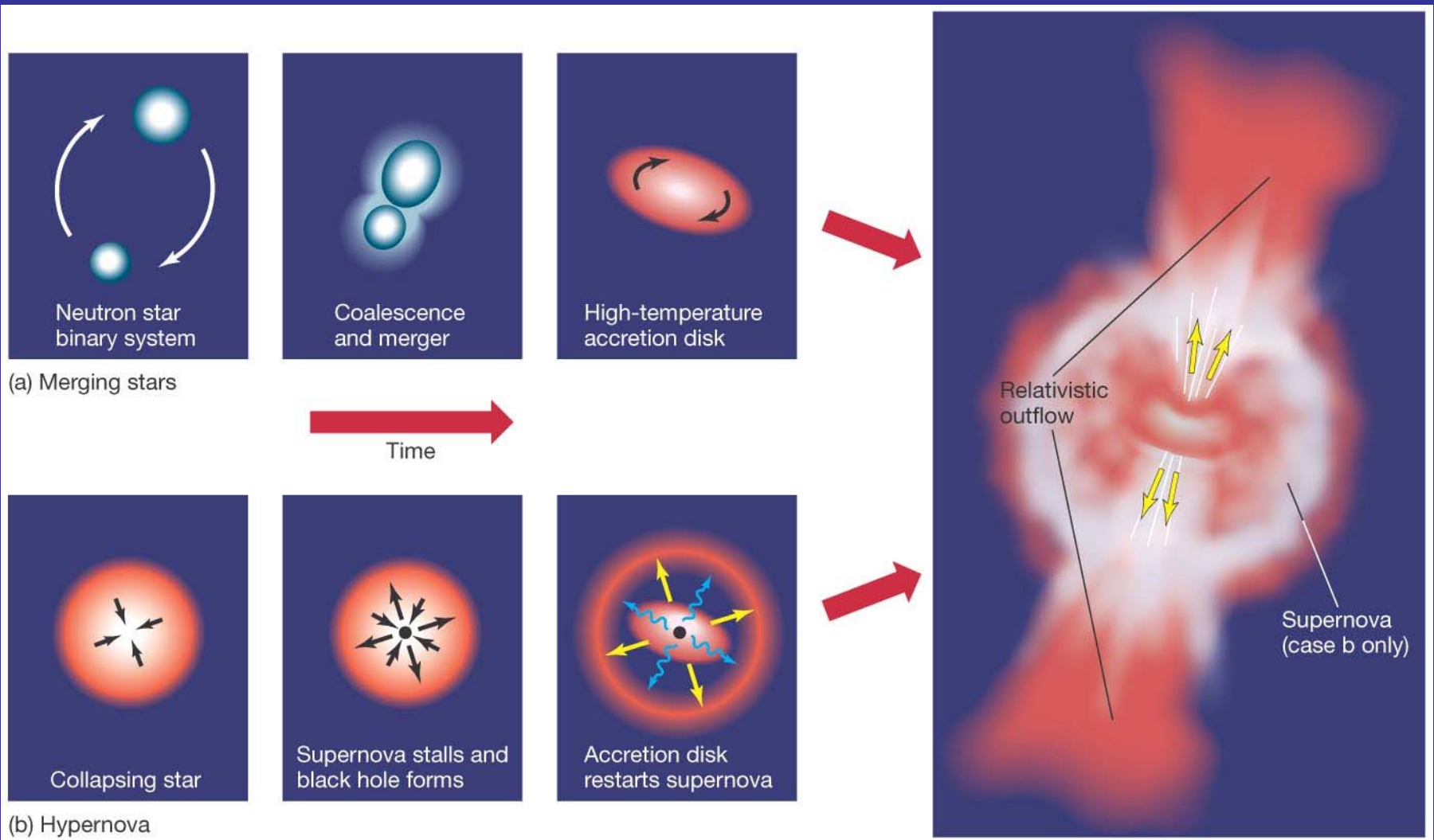
GRB (гамма-всплески): вспышки с энергией от нескольких десятков кэВ до МэВ (иногда и более жесткие). Вспышки длятся от нескольких долей секунд до минут, а иногда и часов.

Короткие гамма-всплески (меньше 2 сек) – слияние нейтронных звезд.

Длинные – Гипернова?

# Gamma-Ray Bursts

Two models—merging Neutron Stars or a “Hypernova” – have been proposed as the source of Gamma-Ray Bursts (“GRB’s”):





# FIRST COSMIC EVENT OBSERVED IN GRAVITATIONAL WAVES AND LIGHT

Colliding Neutron Stars Mark New Beginning of Discoveries

Collision creates light across the entire electromagnetic spectrum. Joint observations independently confirm Einstein's General Theory of Relativity, help measure the age of the Universe, and provide clues to the origins of heavy elements like gold and platinum

Gravitational wave lasted over 100 seconds

On August 17, 2017, 12:41 UTC, LIGO (US) and Virgo (Europe) detect gravitational waves from the merger of two neutron stars, each around 1.5 times the mass of our Sun. This is the first detection of spacetime ripples from neutron stars

Within two seconds, NASA's Fermi Gamma-ray Space Telescope detects a short gamma-ray burst from a region of the sky overlapping the LIGO/Virgo position. Optical telescope observations pinpoint the origin of this signal to NGC 4993, a galaxy located 130 million light years distant



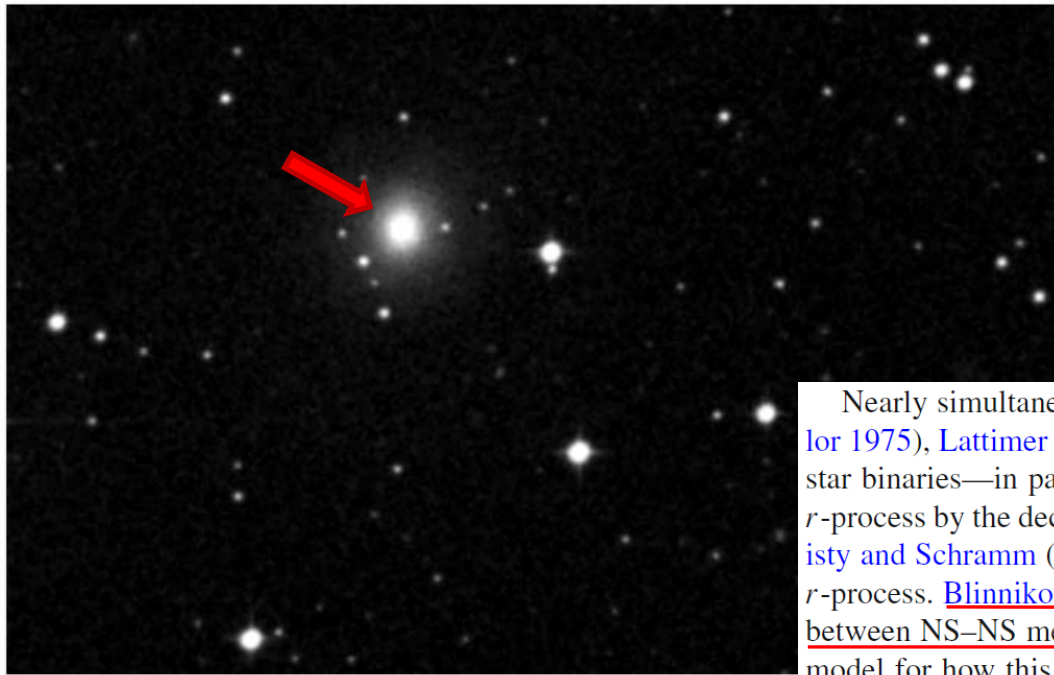
# Rumours swell over new kind of gravitational-wave sighting

Gossip over potential detection of colliding neutron stars has astronomers in a tizzy.

Daive Castelvechi

24 August 2017 | Updated: 25 August 2017, 25 August 2017

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The galaxy NGC 4993 (fuzzy bright spot) in the constellation Hydra, where detectors have spotted gravitational waves from a neutron star merger.

*Событие GW170817*

*Галактика NGC 4993*  
*Расстояние 40 млн пс*



*Гамма-всплеск*  
*GRB 170817A*

Living Rev Relativ (2017) 20:3  
DOI 10.1007/s41114-017-0006-z

REVIEW ARTICLE

## Kilonovae

Brian D. Metzger<sup>1</sup>

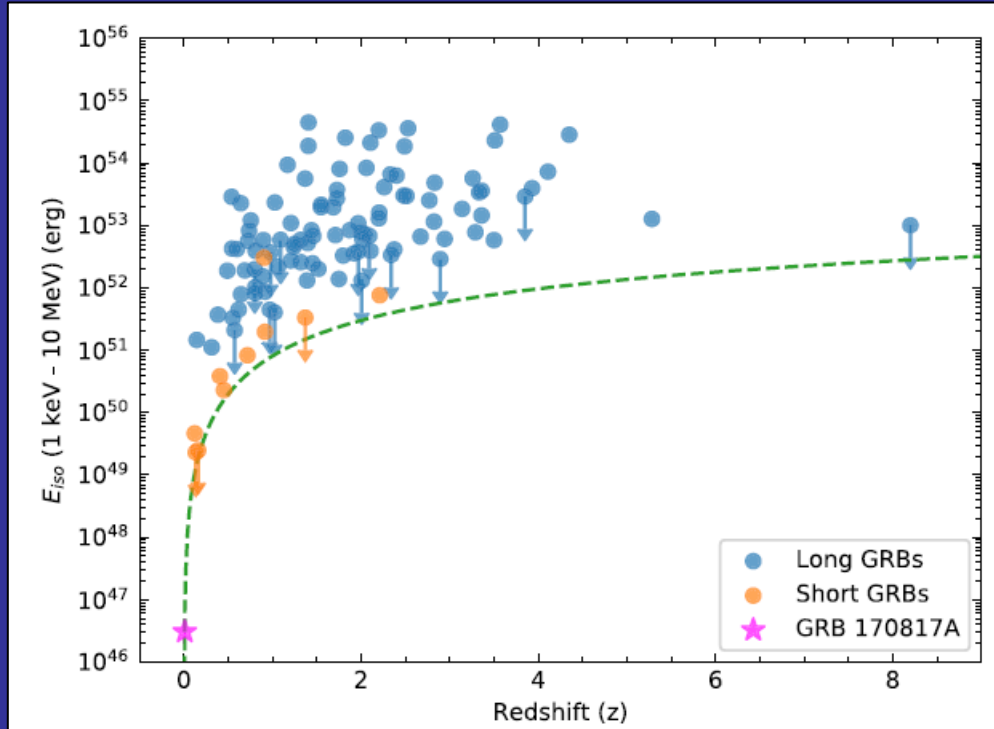
Nearly simultaneous with the discovery of the first binary pulsar (Hulse and Taylor 1975), Lattimer and Schramm (1974, 1976) proposed that the merger of compact star binaries—in particular the collision of BH–NS systems—could give rise to the *r*-process by the decompression of highly neutron-rich ejecta (Meyer 1989). Symbalisty and Schramm (1982) were the first to suggest NS–NS mergers as the site of the *r*-process. Blinnikov et al. (1984) and Paczyński (1986) first suggested a connection between NS–NS mergers and GRBs. Eichler et al. (1989) presented a more detailed model for how this environment could give rise to a GRB (albeit one which differs significantly from the current view). Davies et al. (1994) performed the first numerical simulations of mass ejection from merging neutron stars, finding that ~2% of the binary mass was unbound during the process. Freiburghaus et al. (1999) presented the first explicit calculations showing that the ejecta properties extracted from a hydro-

# Важное о GW170817

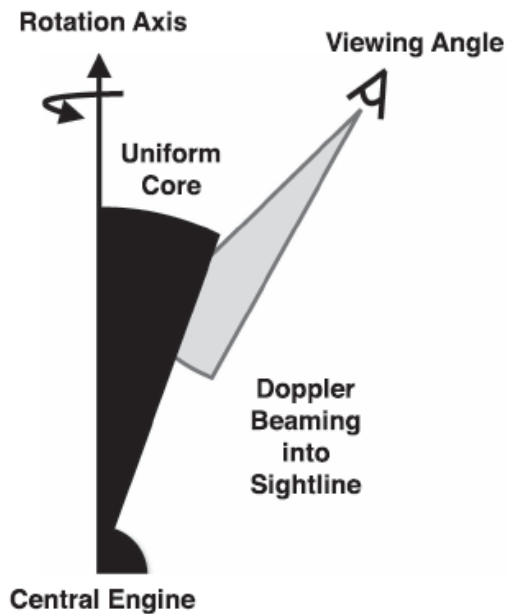
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- ▶ GW170817 – 6-е гравитационно-волновое событие и 1-ое наблюдение слияния объектов с массами нейтронных звезд.
- ▶ Гамма-всплеск GRB170817A наблюдался спустя 1.7 сек. после потери сигнала GW170817.
  - ▶ Подтверждена связь коротких GRB со сливающимися NS
  - ▶ Ограничения на гравитацию: скорость распространения ( $\Delta v/c \lesssim 10^{-15}$ ), лоренц-инвариантность, принцип эквивалентности
- ▶ Спустя 11 часов открыт источник в видимом свете в NGC 4993
  - ▶ Кривые блеска и спектры соответствуют килоновой
  - ▶ Синтез тяжелых элементов в r-процессе
  - ▶ Космология: независимое измерение расстояний, параметра Хаббла
- ▶ Впервые выполнены наблюдения одного объекта в грав.волновом и эл.-маг. (гамма, рентген, ультрафиолет, видимый и инфракрасный свет, радио) канале. Для нейтрино далеко

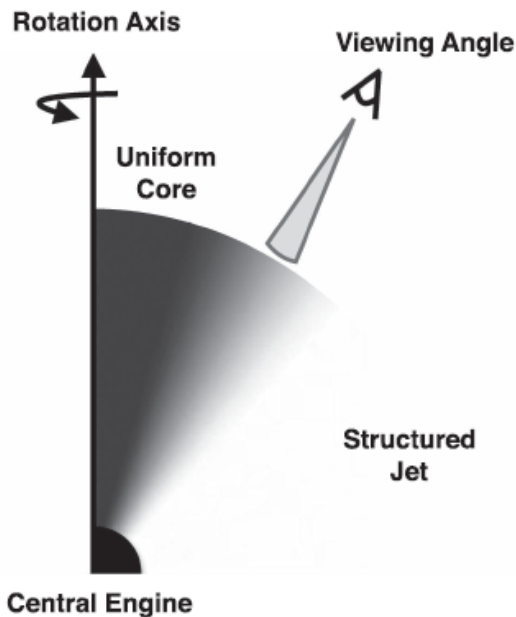
Начало эры многодиапазонной (многоканальной) астрономии – multi-messenger astronomy



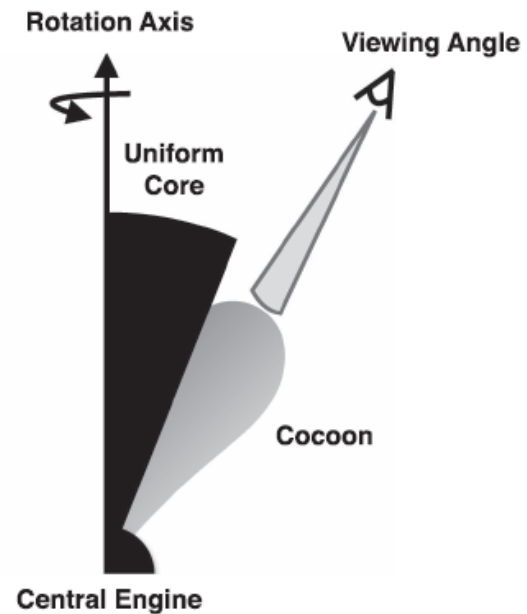
### Scenario i: Uniform Top-hat Jet



### Scenario ii: Structured Jet



### Scenario iii: Uniform Jet + Cocoon





# Exploding neutron stars in close binaries

S. I. Blinnikov, I. D. Novikov, T. V. Perevodchikova, and A. G. Polnarev

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(Submitted January 27, 1984)

*Pis'ma Astron. Zh.* **10**, 422–428 (June 1984)

A close binary system comprising a neutron star and another neutron star (or a black hole) will evolve so that the less massive component sheds mass, passing through a series of quasiequilibrium states, until it achieves its minimum possible mass  $m_{\min} \approx 0.09 M_{\odot}$  and explodes. In a compact globular cluster or the nucleus of a galaxy, such evolution can terminate in an explosion in less than the Hubble time.



# Explosion of a low-mass neutron star

S. I. Blinnikov, V. S. Imshennik, D. K. Nadezhin, I. D. Novikov, T. V. Perevodchikova, and A. G. Polnarev

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(Submitted April 4, 1990)

*Astron. Zh.* **67**, 1181–1194 (November–December 1990)

The process of hydrodynamic destruction of a neutron star that occurs when its mass becomes somewhat less than the minimum mass  $M_{\min} \approx 0.1 M_{\odot}$  is calculated. It is shown that this process occurs explosively and results in the complete dispersal of the neutron star with a kinetic energy  $\sim 4.8$  MeV per nucleon. The calculated results hardly depend on the means by which the mass of the neutron star is reduced to less than  $M_{\min}$  (transfer to a companion in a binary system, decay of nucleons, an equivalent mass decrease due to a decrease in the gravitational constant). Destruction of the neutron star should be accompanied by a short (hundredths of a second) burst of hard thermal x rays and soft gamma rays ( $kT \approx 10$ – $100$  keV), which should be followed by the considerably longer “tail” of x rays and gamma rays associated with a decay of long-lived radioactive nucleons. Some fraction of the explosive energy is carried off in the form of neutrinos.



# Exploding neutron stars in close binaries

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Once having achieved  $m_2 = m_{\min}$ , star 2 will lose its hydrostatic stability and will begin to expand at a rate determined by  $t_{\text{hyd}}$  and the amended equation of state. Clark and Eardley<sup>6</sup> estimate that perhaps one neutron star may undergo tidal disruption every 100 yr within a 15-Mpc radius; thus the event would not be exceedingly rare. Not only should a burst of gravitational waves be produced,<sup>6</sup> but also a powerful electromagnetic flare (most likely x rays and  $\gamma$  rays). Page<sup>2</sup> believes that the explosion may attain an energy of supernova scale, but the problem awaits a detailed analysis. We intend to consider this process further in a separate paper.

We also have omitted discussion here of the physical processes that will accompany the mass transfer, such as the stripping from the star of material with nuclei having excess neutrons; as these nuclei later decay,  $\gamma$ -ray burster phenomena might occur (like the processes that Bisnovatyi-Kogan and Checkëtkin<sup>13</sup> have discussed).

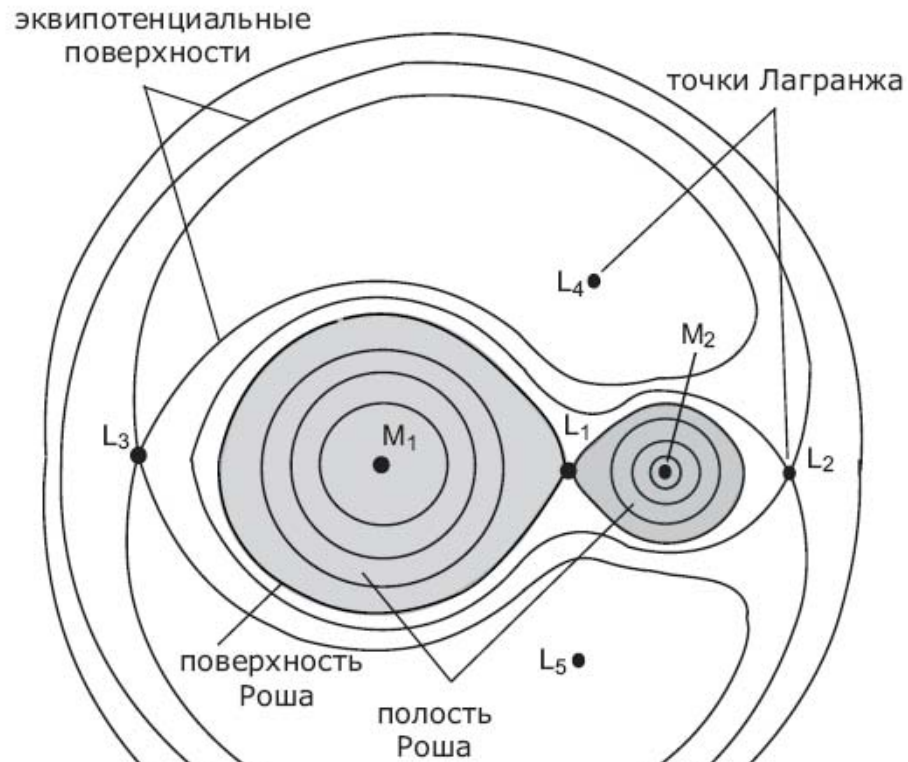
$$t_{\text{grav}} \approx (10^{10} \text{ yr}) \left( \frac{m_1}{10^{33} \text{ g}} \right)^{-1} \left( \frac{m_2}{10^{33} \text{ g}} \right)^{-1} \left( \frac{m_1}{10^{33} \text{ g}} + \frac{m_2}{10^{33} \text{ g}} \right)^{-1} \left( \frac{a}{R_{\odot}} \right)^4$$

$$\Phi = -\frac{GM_1}{r_1} - \frac{GM_2}{r_2} - \frac{1}{2}\omega^2[(x - \mu a)^2 + y^2]$$

$$\omega = \sqrt{\frac{G(M_1 + M_2)}{a^3}}$$

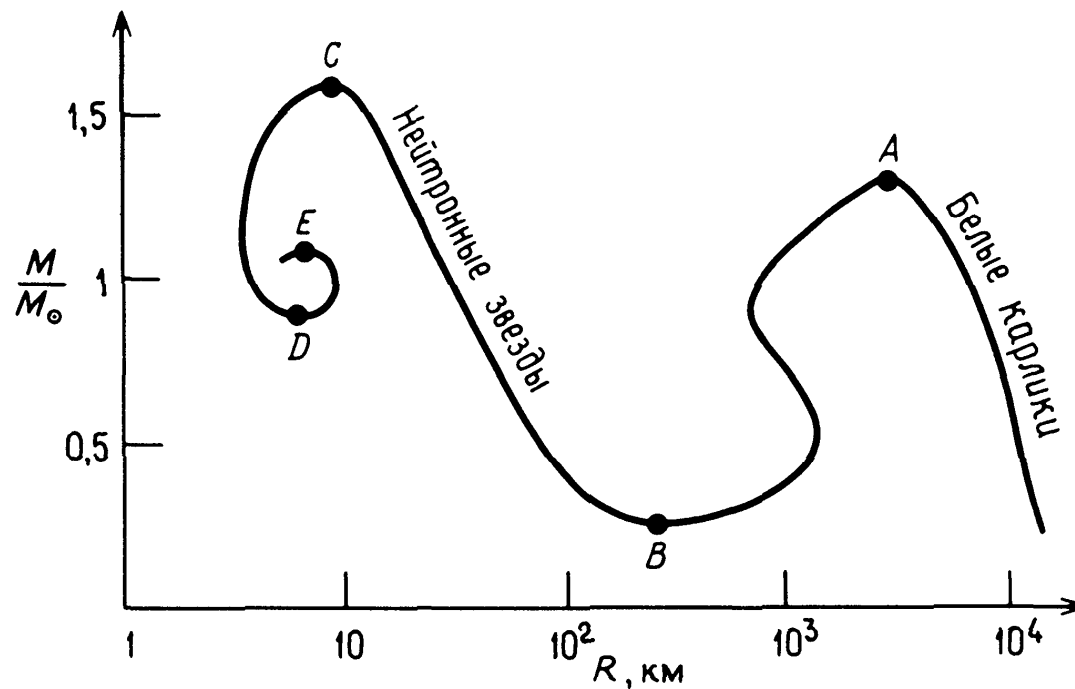
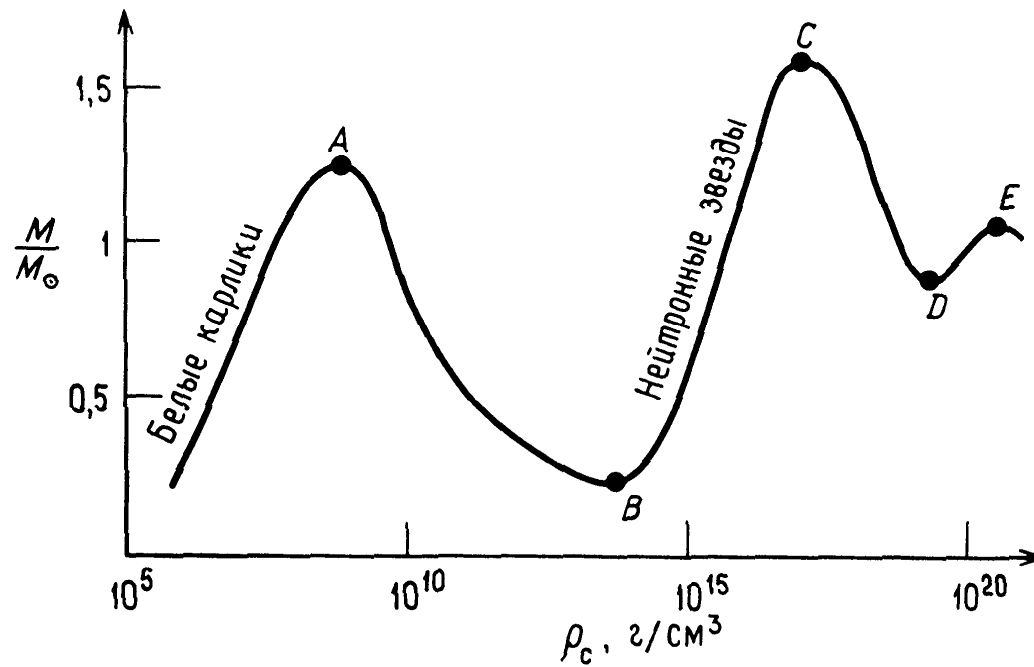
$$L = \sqrt{G \frac{M_1^2 M_2^2}{M_1 + M_2} a}$$

$$\frac{d \ln R_2}{d \ln M_2} > 2 \frac{M_2}{M_1} - \frac{5}{3}$$



Stuart L. Shapiro  
Saul A. Teukolsky

# Black Holes, White Dwarfs, and Neutron Stars





## EVOLUTION OF CLOSE NEUTRON STAR BINARIES

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Observatory and Department of Physics, Yale University

*Received 1976 November 11; revised 1976 December 17*

### ABSTRACT

In binary systems consisting of two neutron stars, the orbit decays by gravitational radiation. A crude model shows that the less massive star may suffer either immediate tidal disruption or slow mass stripping when it reaches its Roche radius, depending on the initial masses and on the details of mass exchange or mass loss. Typical energy releases are  $4 \times 10^{52}$  ergs in gravitational waves before the onset of stripping,  $2 \times 10^{52}$  ergs in gravitational waves after the onset of stripping,  $2 \times 10^{53}$  ergs in neutrinos after the onset of stripping. The stripping process always ends in tidal disruption of the less massive star after a few seconds or a few hundred revolutions.

As the endpoint of binary stellar evolution, such events are estimated to occur only every  $\sim 100$  yr out to a radius of 15 Mpc, and are thus less important than supernovae as sources of gravitational waves; the observed wave amplitude would be  $h \sim 10^{-21}$ . Such events may occur in Type II supernovae, if the collapsing stellar core rotates rapidly enough to fission into two neutron stars.

*Subject headings:* gravitation — stars: binaries — stars: evolution — stars: neutron

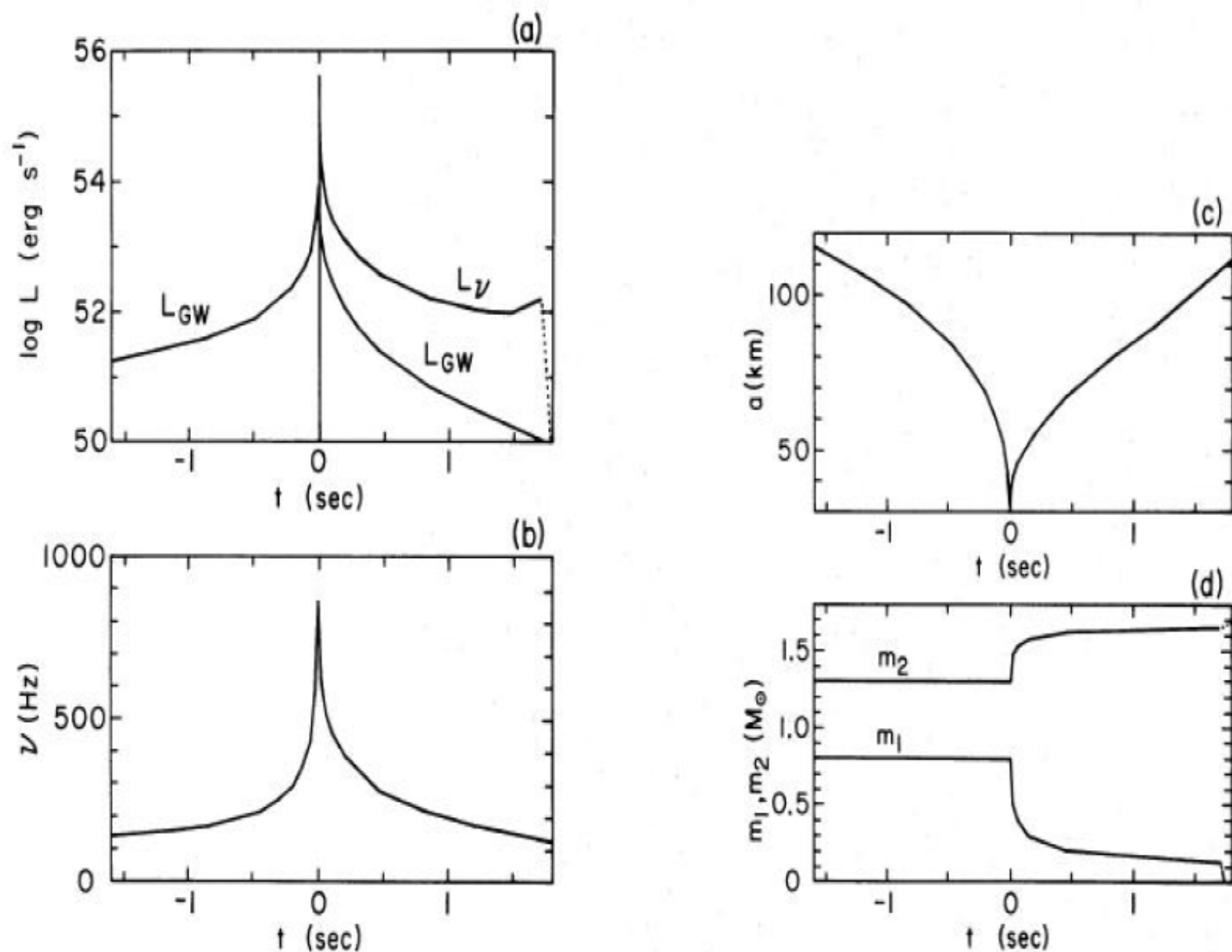
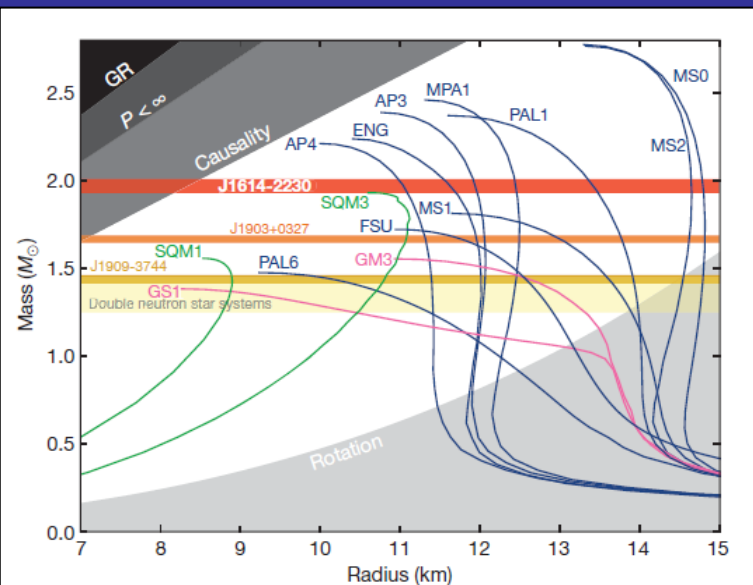
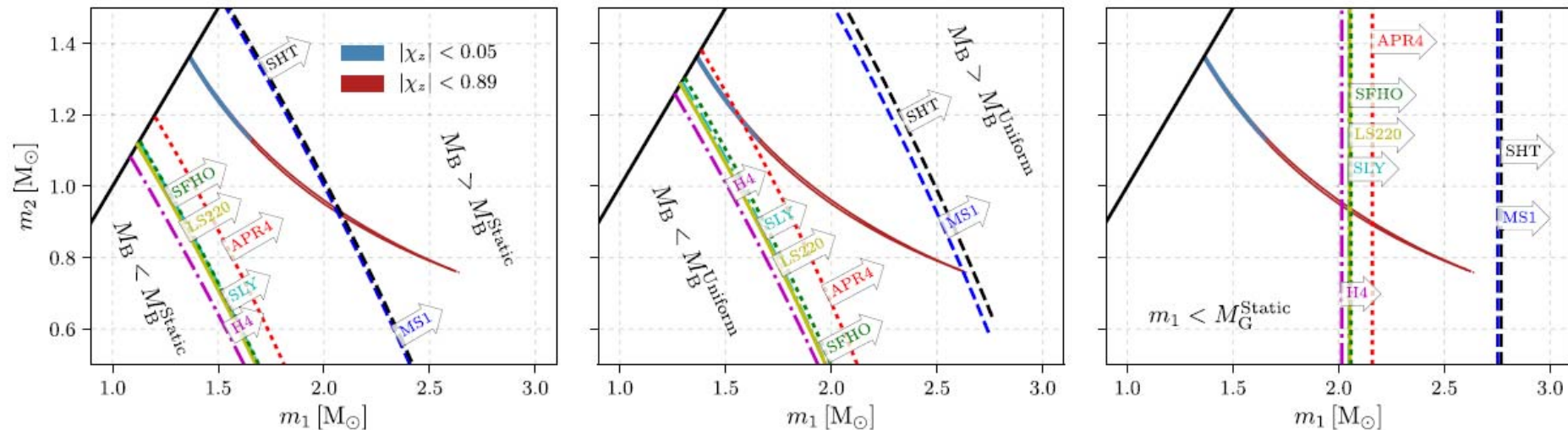


FIG. 7—Time evolution of a system with initial masses  $0.8$  and  $1.3 M_{\odot}$ . (a) Neutrino and gravitational wave luminosities. (b) Frequency of gravitational wave. (c) Separation of components. (d) Masses of stars.

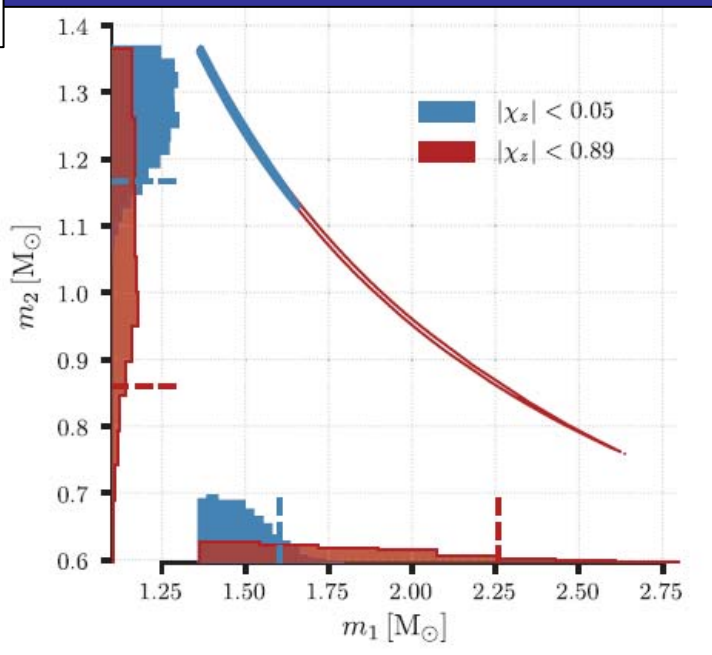


$$M = \frac{(m_1 m_2)^{3/5}}{(m_1 + m_2)^{1/5}}$$

← *Chirp mass*

**high-spin:**  
 $m_1 \in (1.36 \div 2.26) M_\odot$   
 $m_2 \in (0.86 \div 1.36) M_\odot$   
 $m_{tot} = 2.82^{+0.47}_{-0.09} M_\odot$

**low-spin:**  
 $m_1 \in (1.36 \div 1.60) M_\odot$   
 $m_2 \in (1.17 \div 1.36) M_\odot$   
 $m_{tot} = 2.74^{+0.04}_{-0.01} M_\odot$





# INTEGRAL and NASA's Fermi satellite

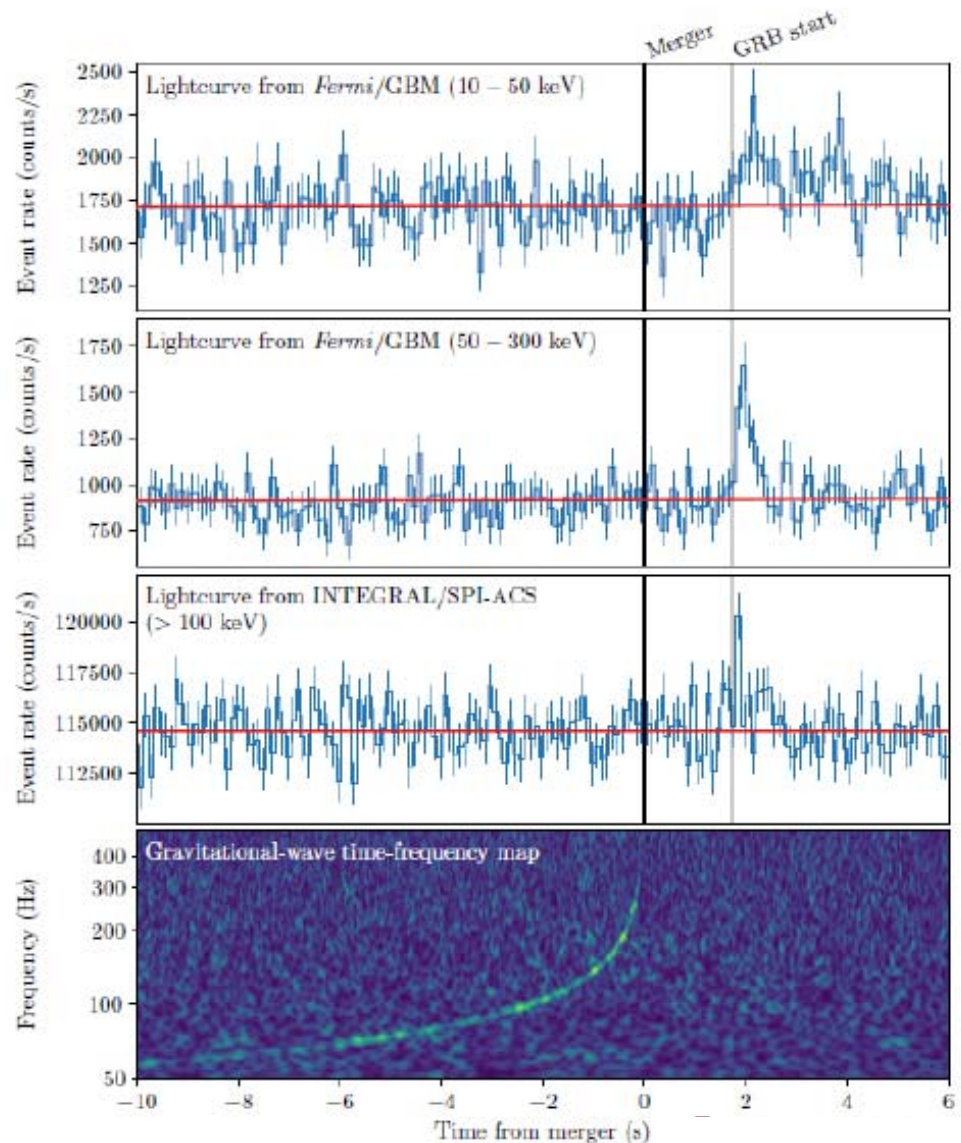


(INTEGRAL)

INTERNATIONAL  
Gamma-Ray  
Astrophysics  
Laboratory



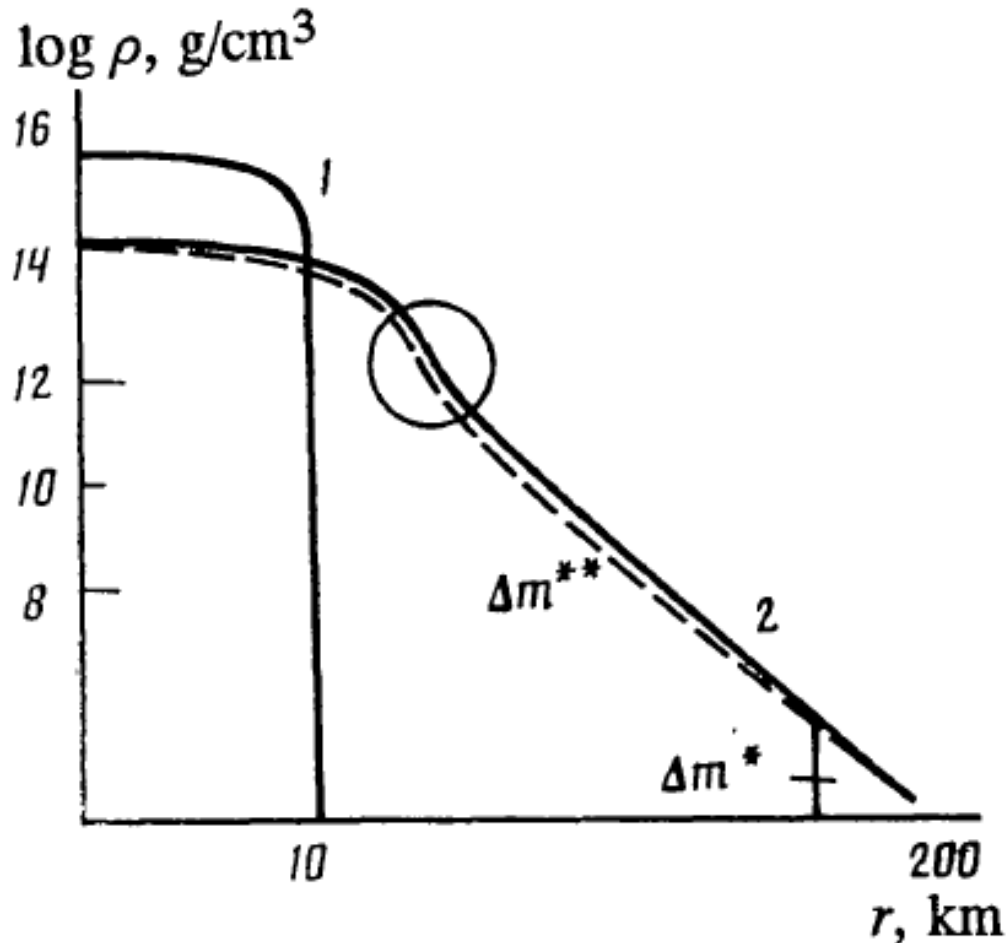
The Fermi  
Gamma-ray  
Space  
Telescope



# Постановка задачи о взрыве

## Explosion of a low-mass neutron star

S. I. Blinnikov, V. S. Imshennik, D. K. Nadezhin, I. D. Novikov, T. V. Perevodchikova, and A. G. Polnarev

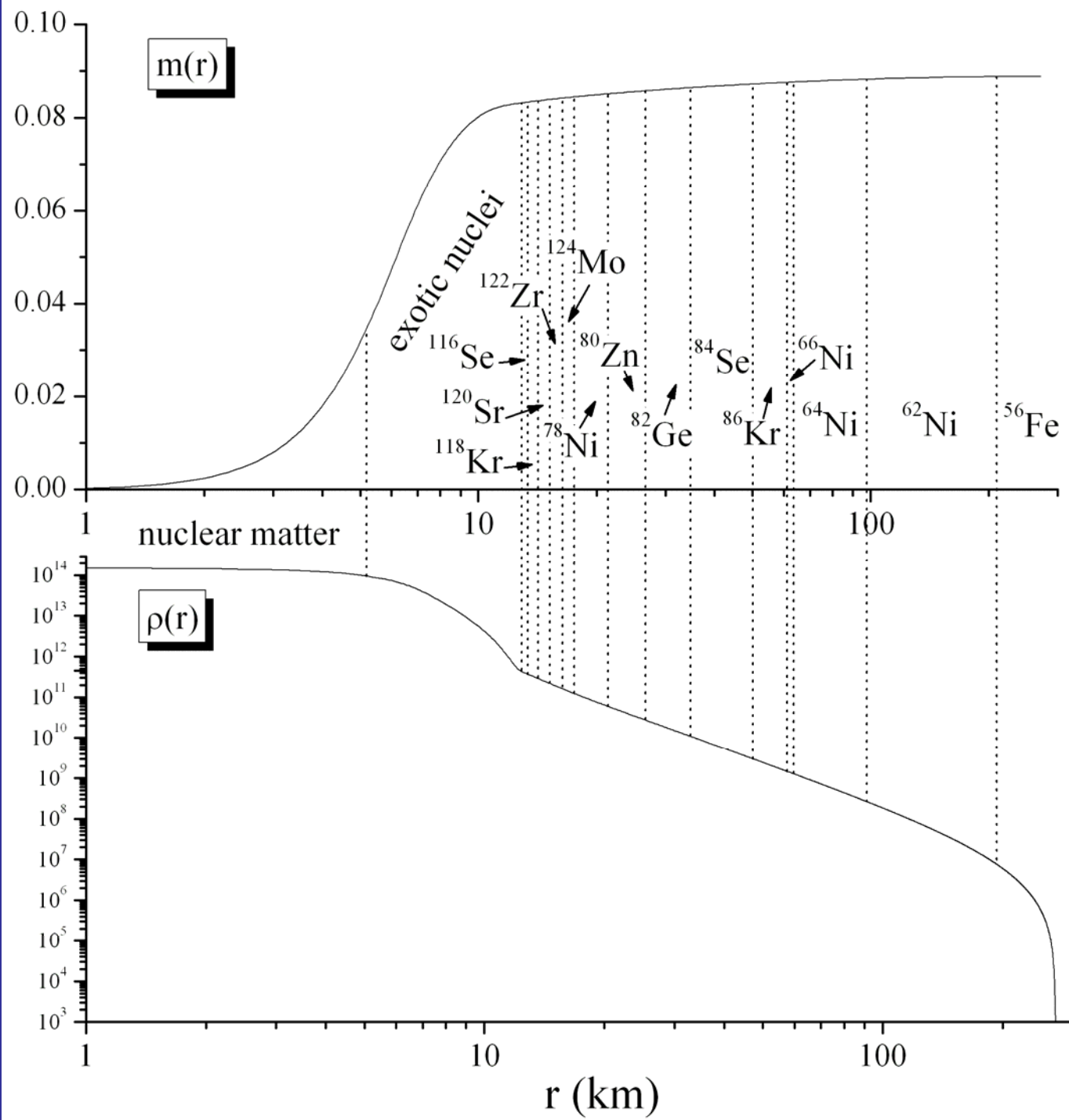


$$P = P_0(\rho) + \rho \tilde{R}T + \frac{1}{3} aT^4,$$

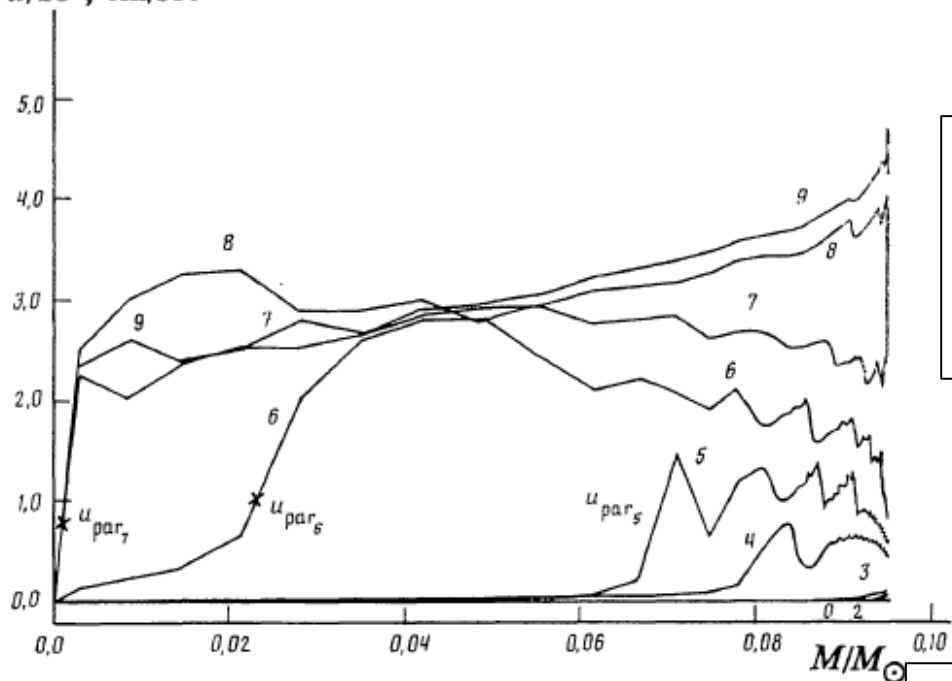
$$E = E_0(\rho) + \frac{3}{2} \tilde{R}T + \frac{aT^4}{\rho}.$$

$$t_{\text{hyd}} \approx \frac{1}{\sqrt{6\pi G \bar{\rho}}} \approx 0.3 \text{ msec}$$

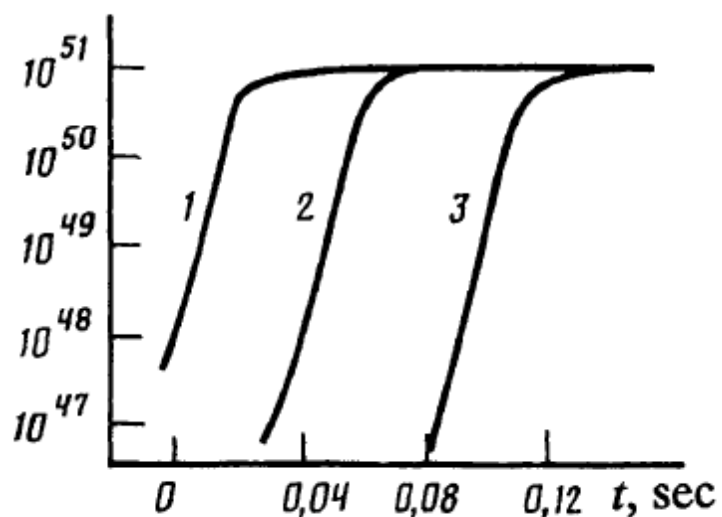
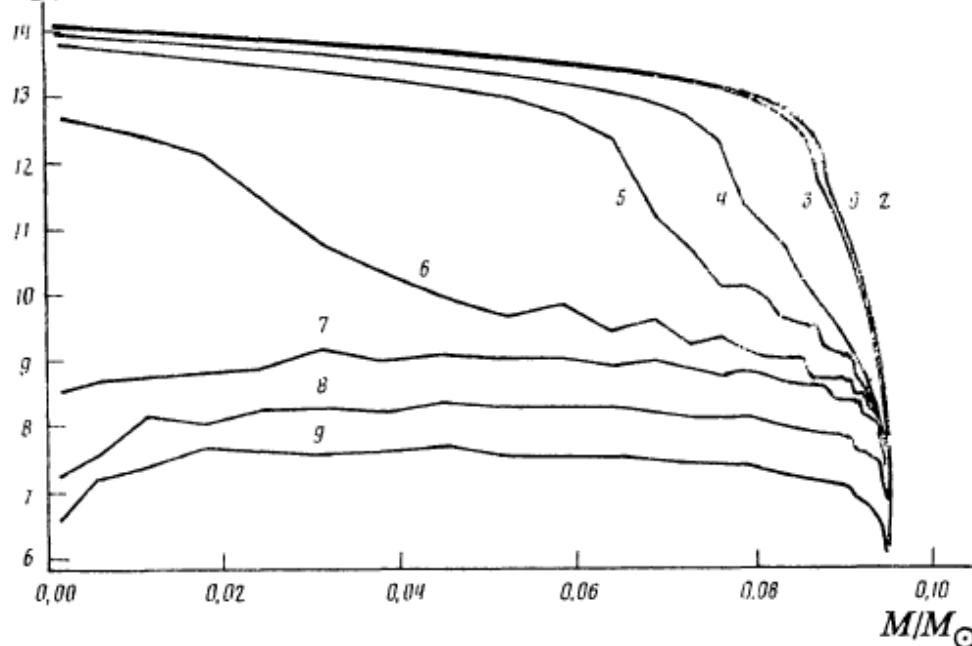
$$\left\{ \begin{array}{l} \frac{\partial r}{\partial t} = u, \\ \frac{\partial u}{\partial t} = -4\pi r^2 \frac{\partial P}{\partial m} - \frac{Gm}{r^2}, \\ \frac{1}{\rho} = \frac{4\pi}{3} \frac{\partial r^3}{\partial m}, \\ \frac{\partial E}{\partial t} + P \frac{\partial 1/\rho}{\partial t} = 0. \end{array} \right.$$





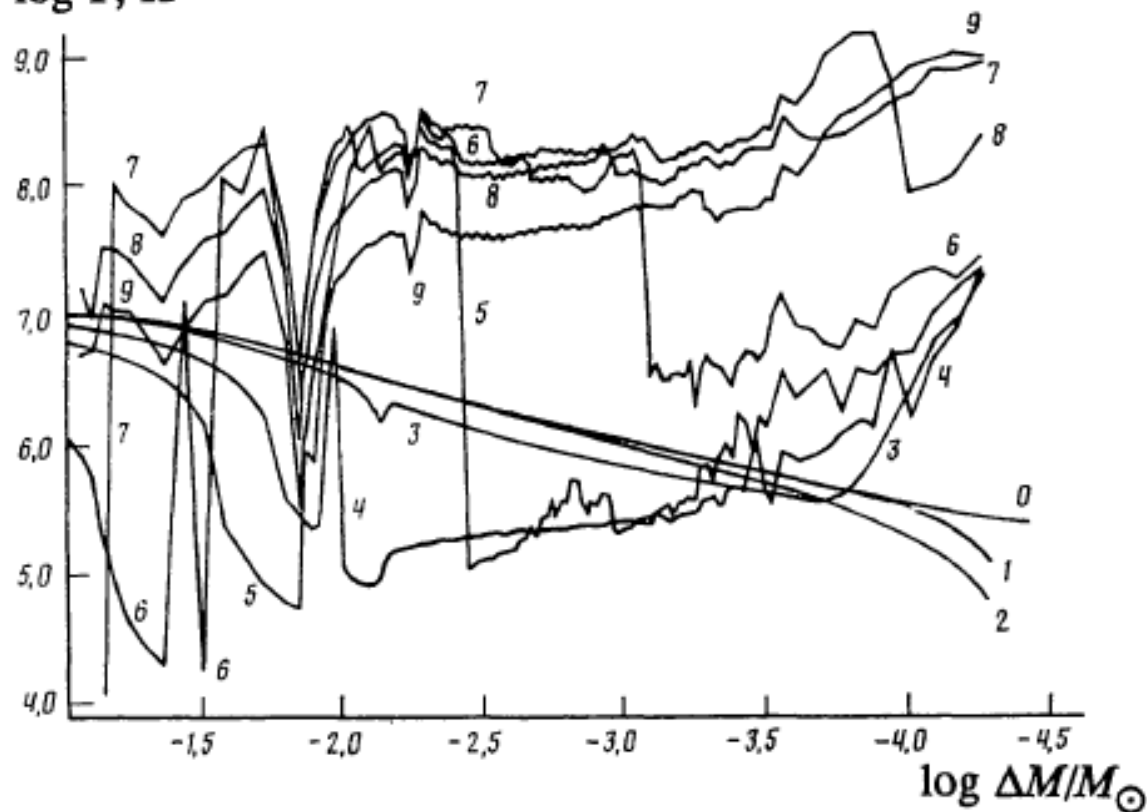
$u/10^9, \text{ cm/sec}$ 

| Curve No. in the figures | Curve No. in the figures | Curve No. in the figures | Time, sec            | Curve No. in the figures | Time, sec            |
|--------------------------|--------------------------|--------------------------|----------------------|--------------------------|----------------------|
| 0                        | 0                        | 4                        | $9,70 \cdot 10^{-2}$ | 7                        | $1,20 \cdot 10^{-1}$ |
| 1                        | $5,60 \cdot 10^{-3}$     | 5                        | $1,03 \cdot 10^{-1}$ | 8                        | $1,31 \cdot 10^{-1}$ |
| 2                        | $2,63 \cdot 10^{-2}$     | 6                        | $1,11 \cdot 10^{-1}$ | 9                        | $1,44 \cdot 10^{-1}$ |
| 3                        | $5,52 \cdot 10^{-2}$     |                          |                      |                          |                      |

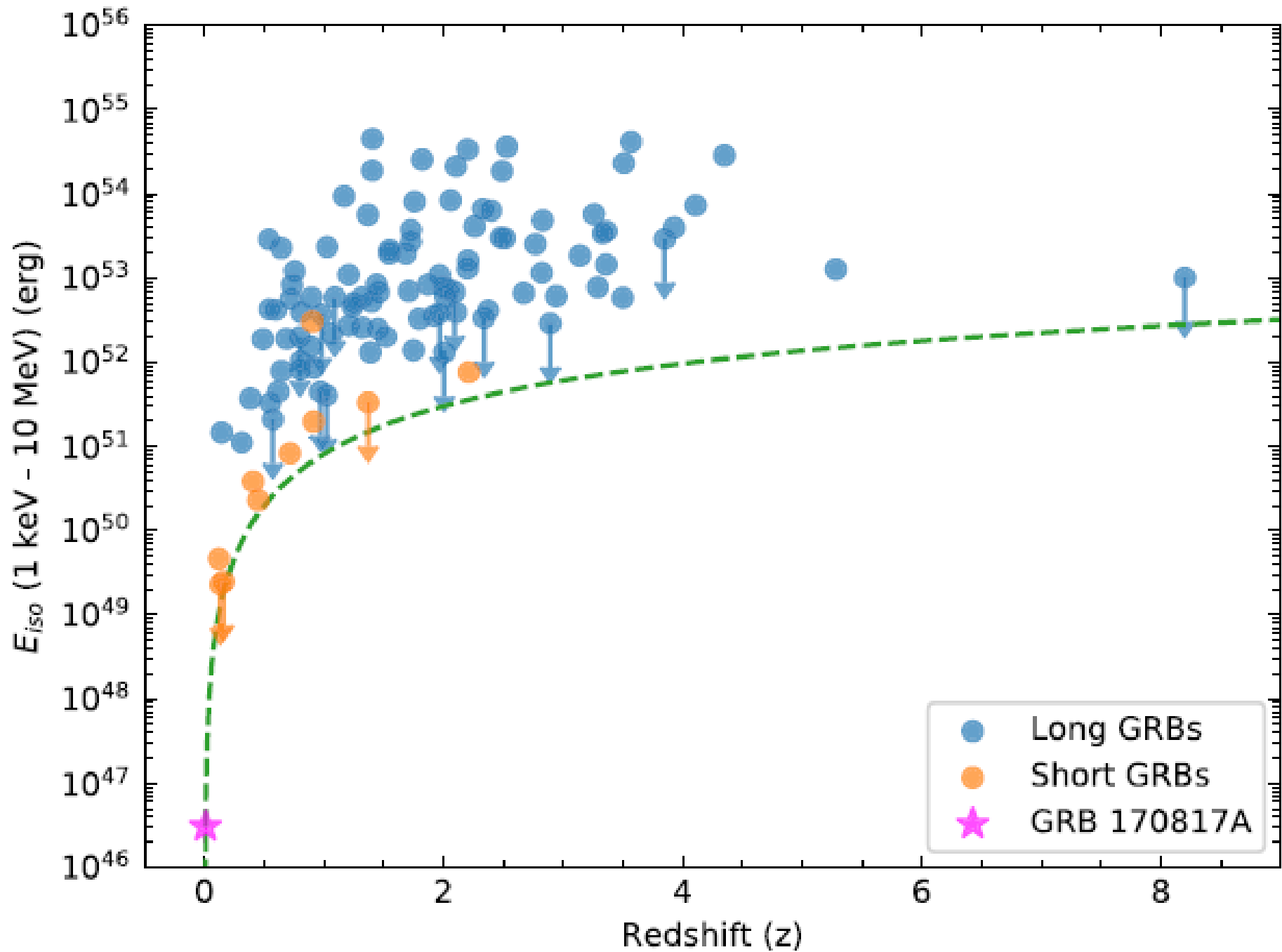
 $\mathcal{E}_{\text{kin}}, \text{ erg}$  $\log \rho, \text{ g/cm}^3$ 

surface layers can be maintained at  $10^8$ - $10^9$  K. This should lead to a burst of hard thermal x rays and soft gamma rays with a total energy of  $10^{43}$ - $10^{47}$  erg.

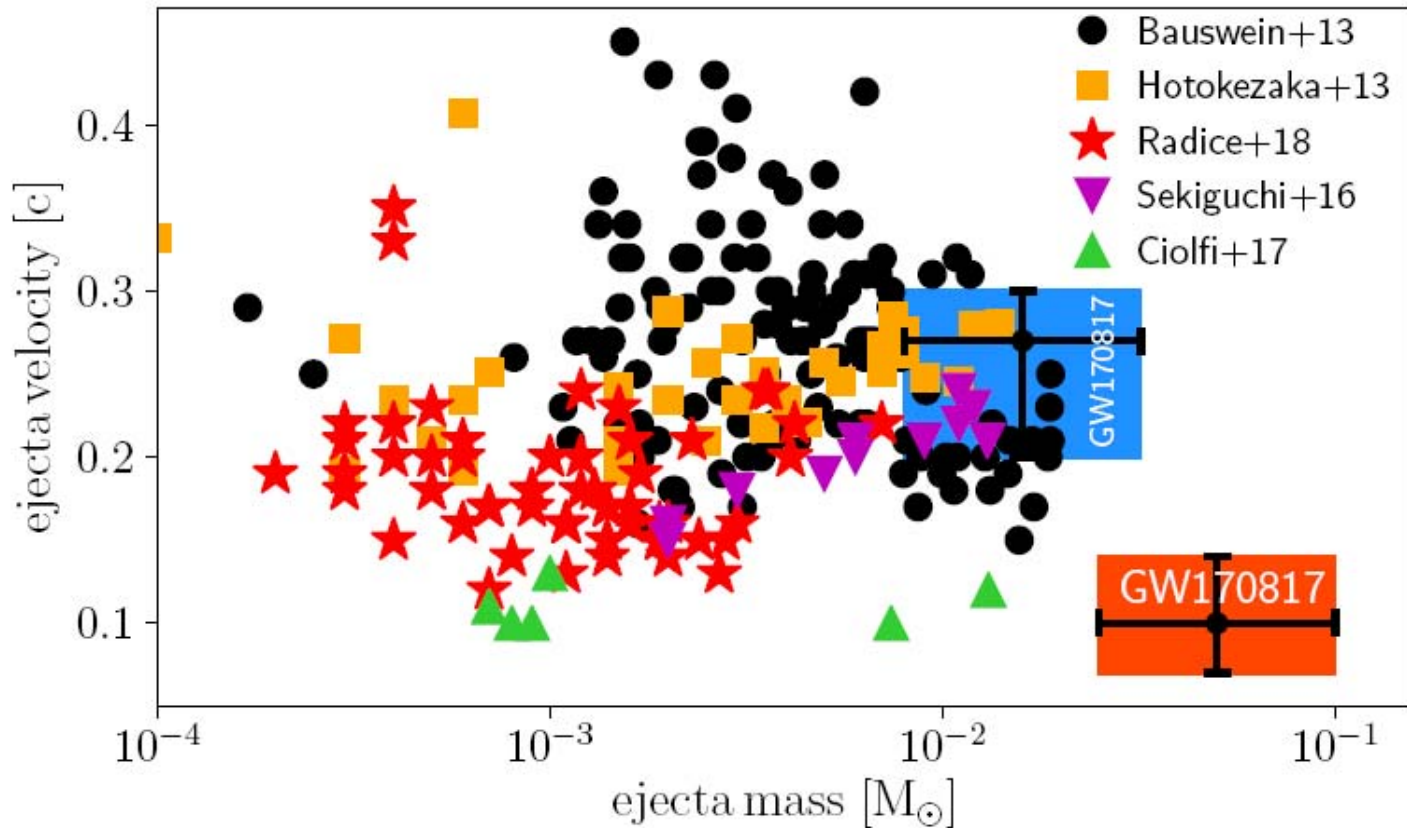
log T, K



**FIG. 12.** Temperature distributions along the Lagrangian coordinate ( $\Delta M$  is the mass reckoned from the surface) in the course of the explosion of a neutron star of mass  $M = 0.09499 M_{\odot}$  at different times (Table I). The temperature increase to  $10^8$ - $10^9$  K at the surface indicates the possibility of thermal x-ray and gamma-ray bursts accompanying the explosions of neutron stars.



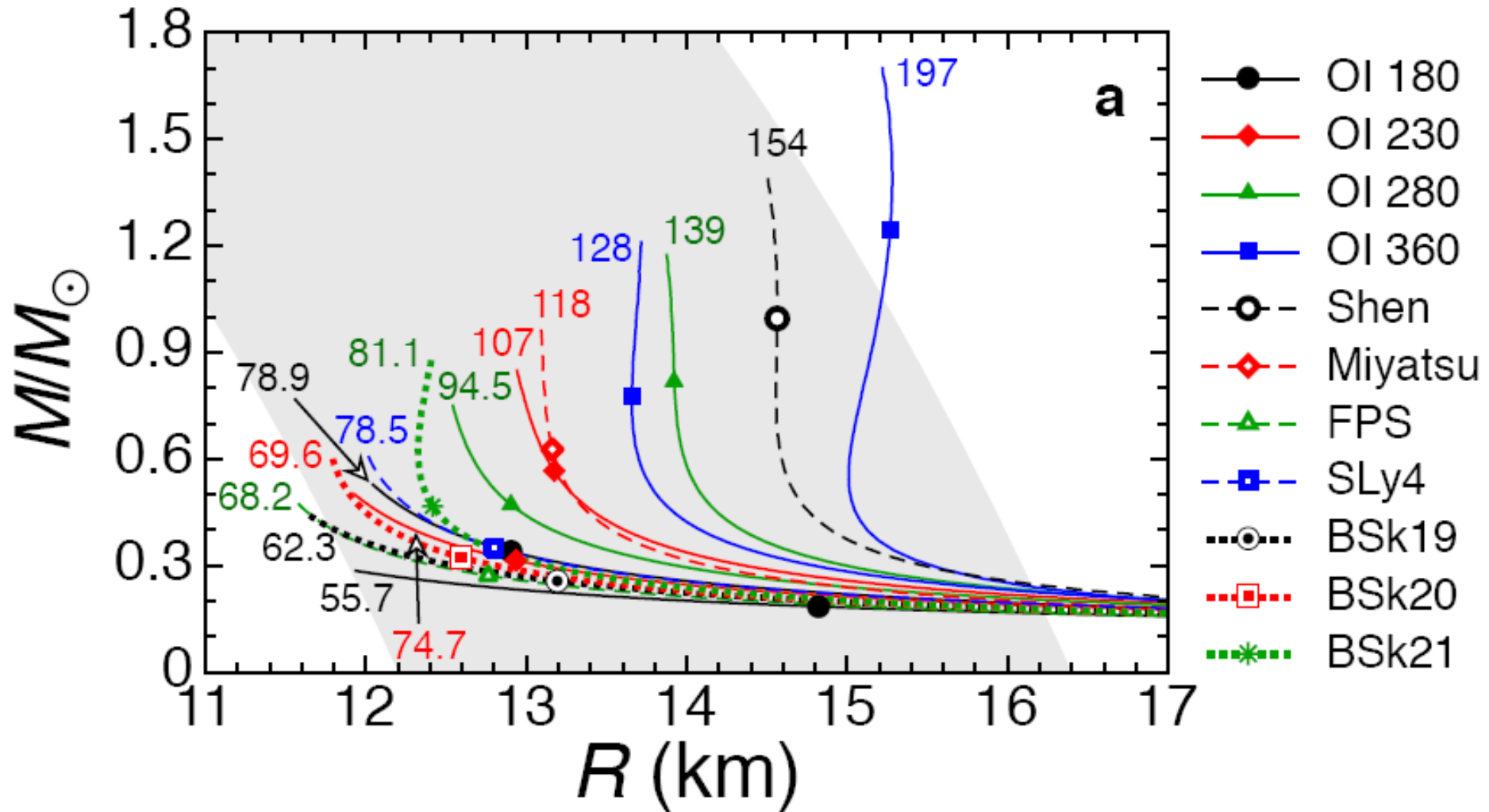


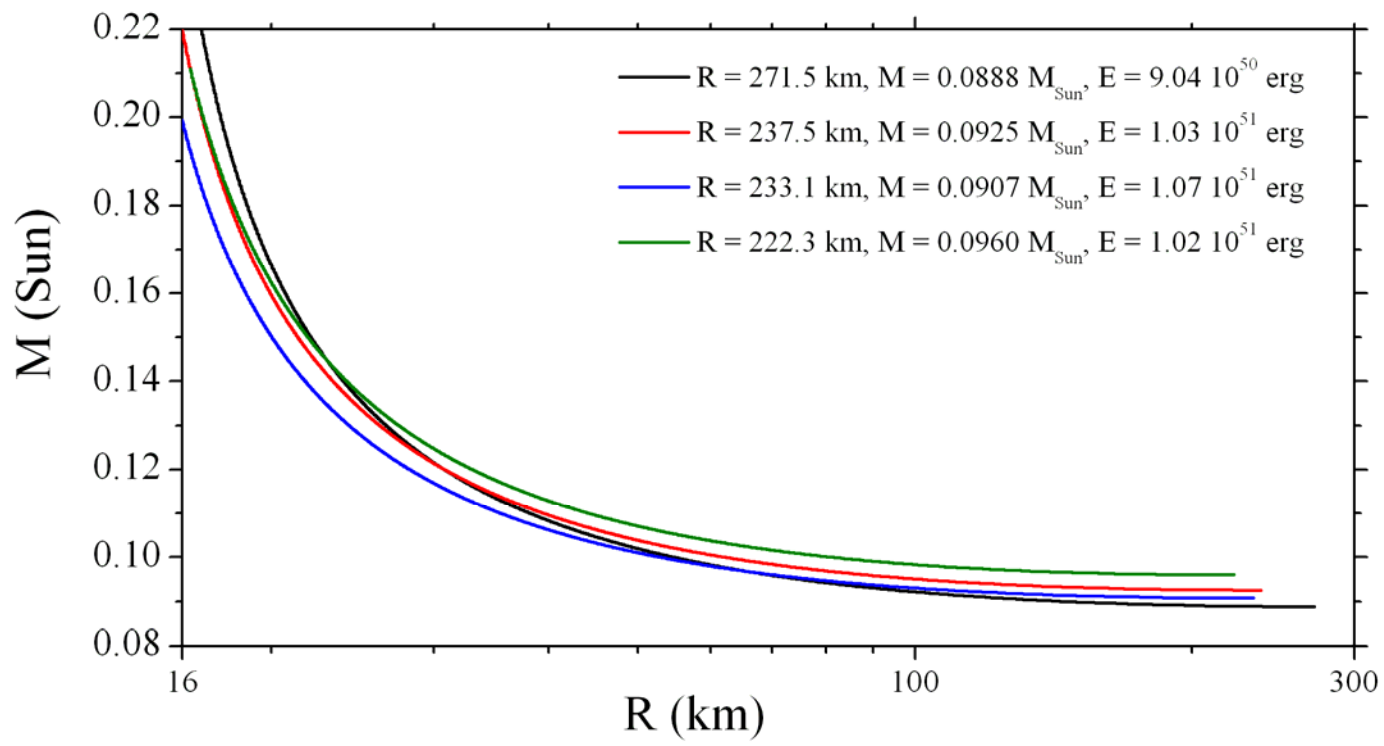
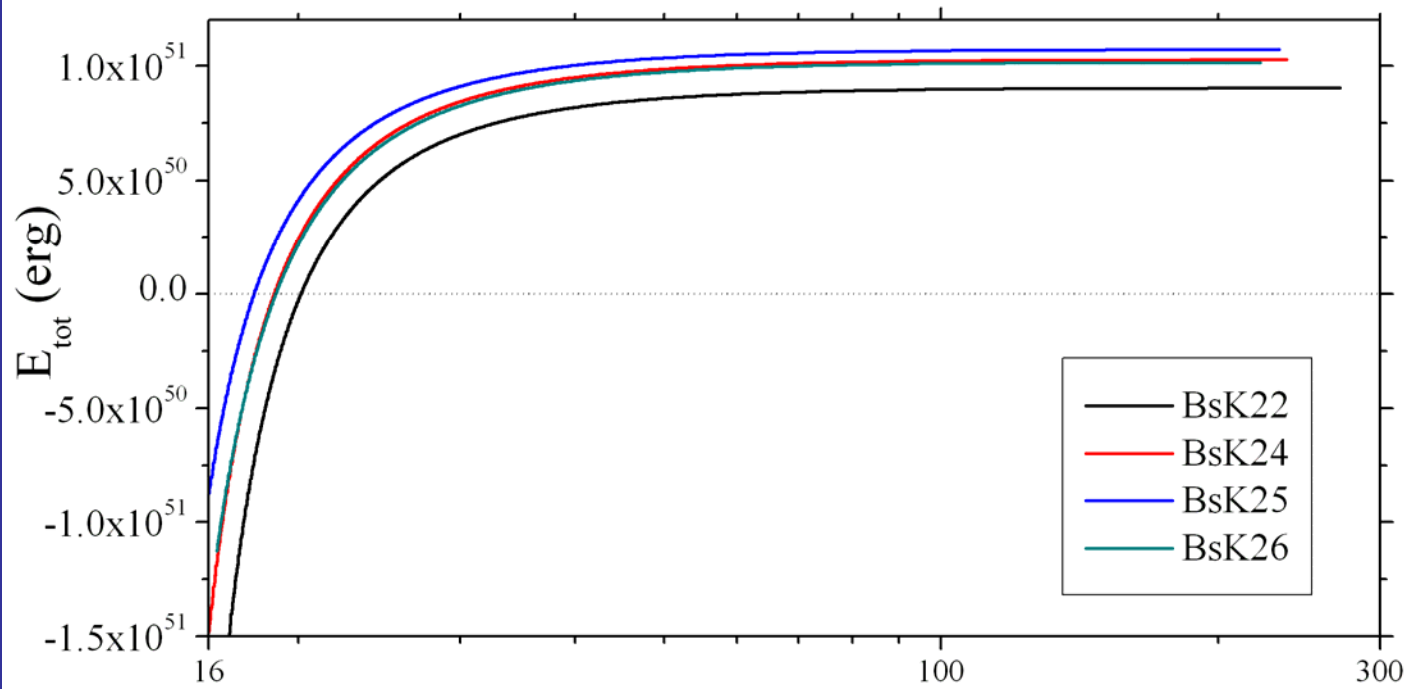


**Figure 1.** Dynamical ejecta masses and velocities from various binary neutron star merger simulations encompassing different numerical techniques, various equations of state, binary binary mass ratios 0.65 – 1.0, effects of neutrinos and magnetic fields [77,78,64,73,65], together with the corresponding ejecta parameters inferred from the ‘blue’ and ‘red’ kilonova of GW170817 (see the text for details).

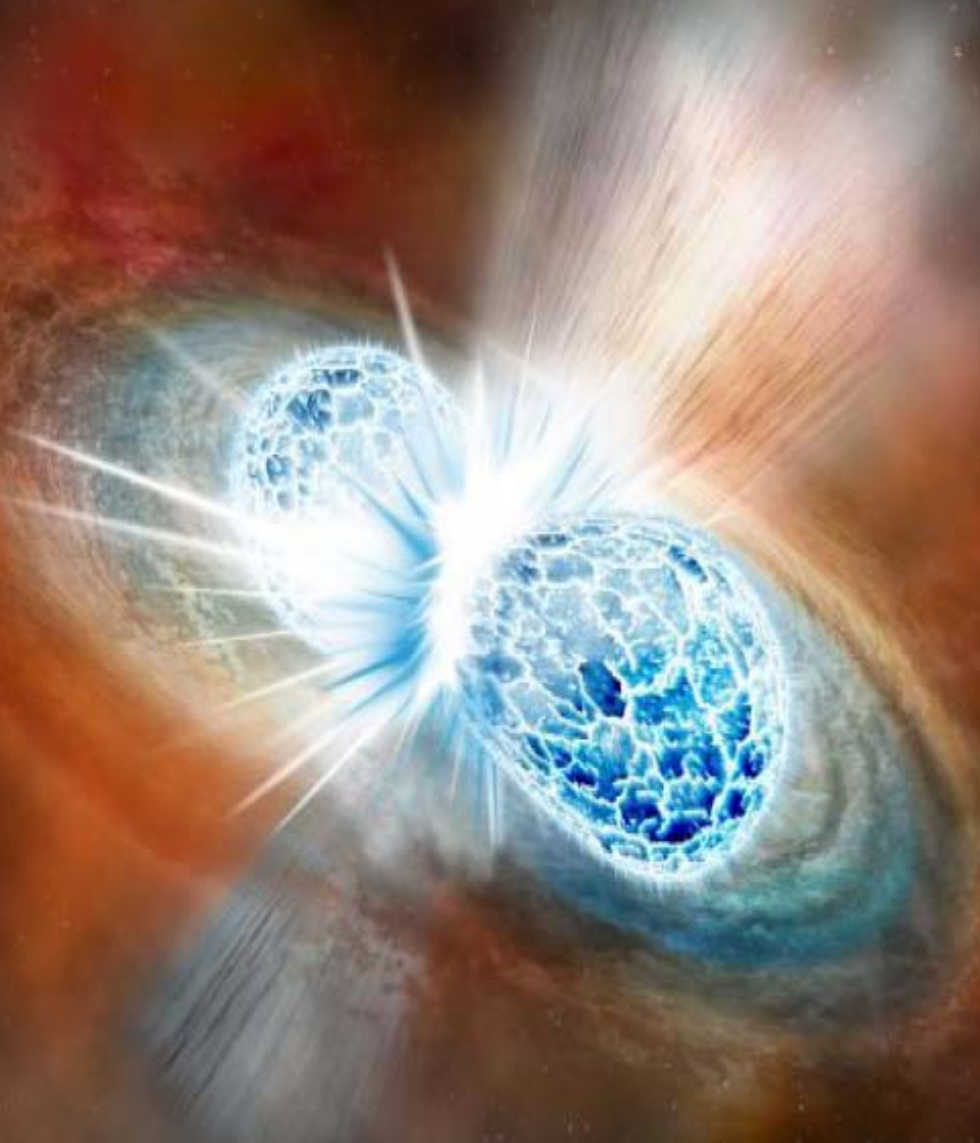
# Mass and radius formulas for low-mass neutron stars

Hajime Sotani<sup>1,\*</sup>, Kei Iida<sup>2</sup>, Kazuhiro Oyamatsu<sup>3</sup>, and Akira Ohnishi<sup>1</sup>





# Спасибо за внимание!



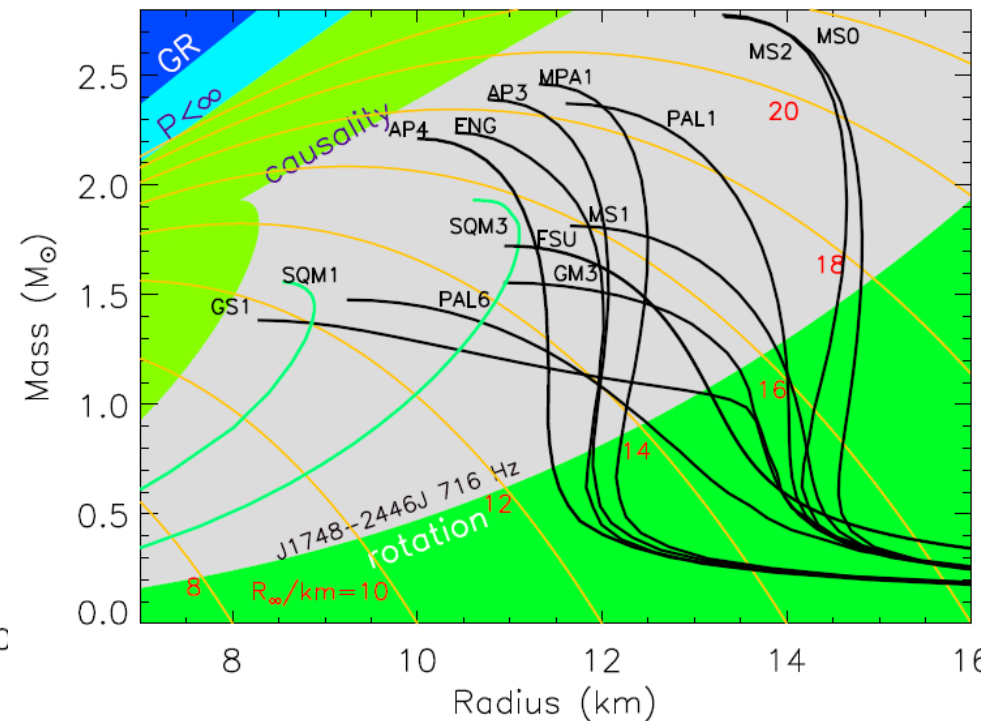
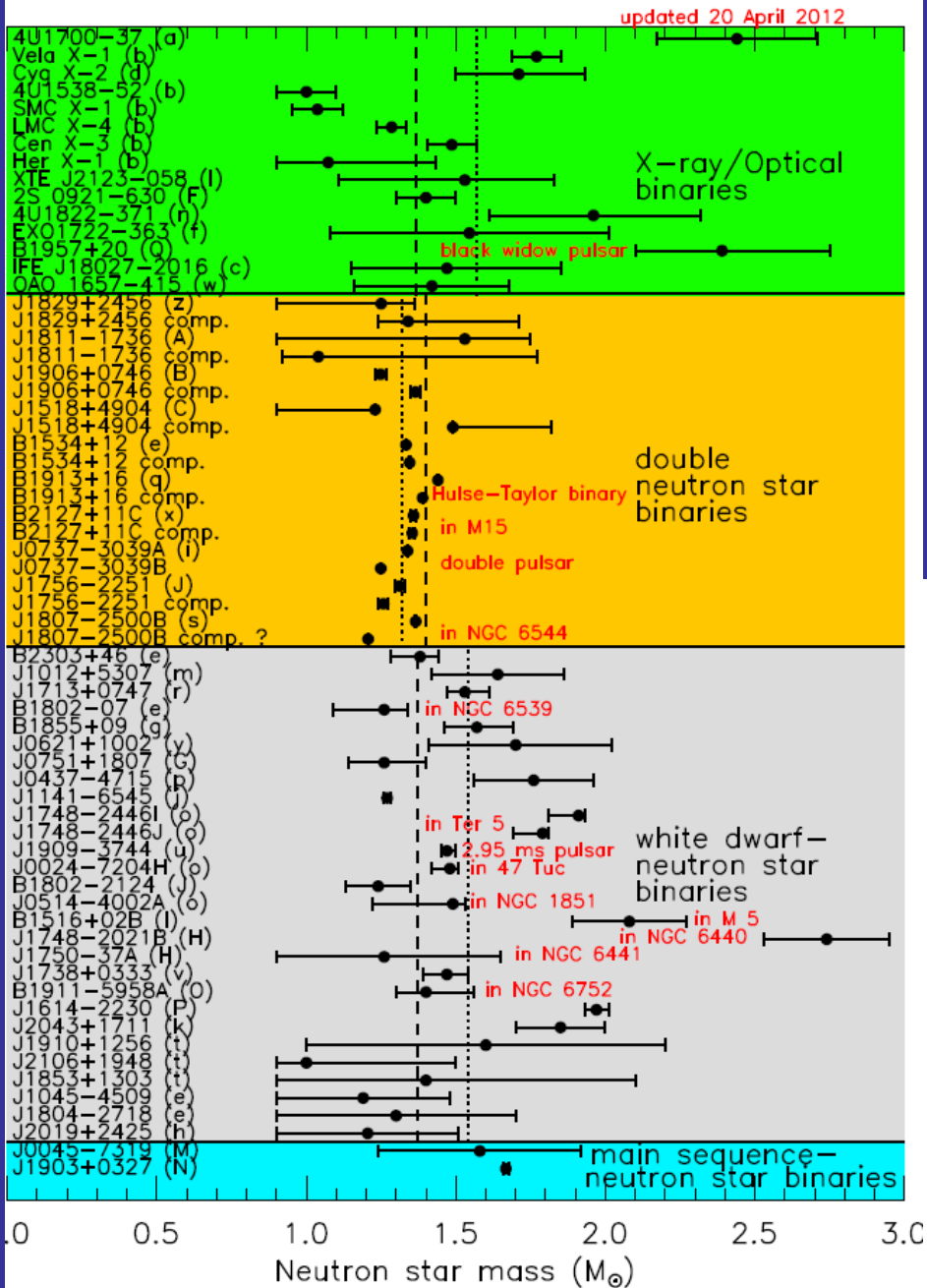
Caption: Artist's concept of the explosive collision of two neutron stars. Illustration by Robin Dienel courtesy of the Carnegie

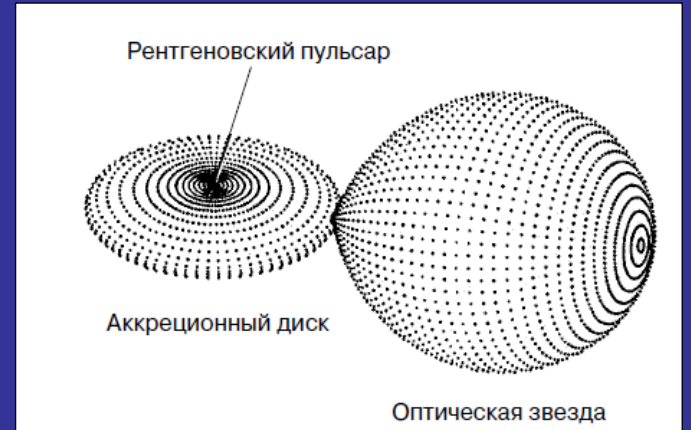
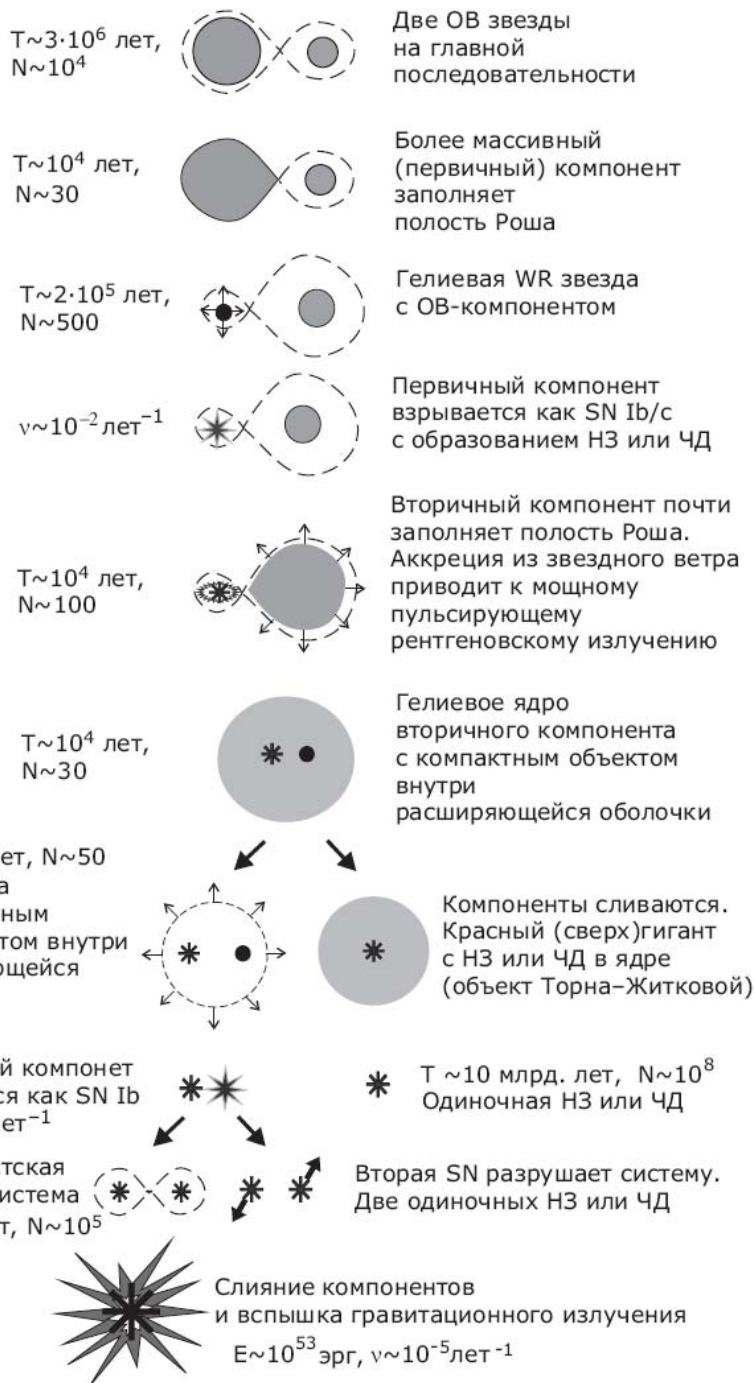


# Maximum neutron star mass

J.M. Lattimer

Annual Review of Nuclear and Particle Science, vol. 62, issue 1, pp. 485-515 (2012)





А.В. Тутуков и Л.Р. Юнгельсон,  
 Научные информ. Астросовета АН  
 СССР. Т. 27. 1973

# M-R diagram for neutron stars

